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Wenjun Dong Embry-Riddle Aeronautical University, dongw1@erau.edu

David Fritts GATS Boulder

Alan Z. Liu Embry Riddle Aeronautical University - Daytona Beach, liuz2@erau.edu

Hanli Liu National Center for Atmospheric Research

Jonathan Snively Embry-Riddle Aeronautical University, snivelyj@erau.edu

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# Accelerating Atmospheric Gravity Wave Simulations using Machine Learning: Kelvin-Helmholtz Instability and Mountain Wave Sources Driving Gravity Wave Breaking and Secondary Gravity Wave Generation

## Wenjun DONG

Embry-Riddle Aeronautical University https://orcid.org/0000-0003-0773-5862

## **David Fritts**

GATS Boulder

## **Thomas Lund**

GATS Boulder

## Alan Liu ( liuz2@erau.edu)

Embry-Riddle Aeronautical University https://orcid.org/0000-0002-1834-7120

## Hanli Liu

National Center for Atmospheric Research https://orcid.org/0000-0002-6370-0704

## **Jonathan Snively**

Embry-Riddle Aeronautical University https://orcid.org/0000-0002-7616-439X

## Article

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1	Accelerating Atmospheric Gravity Wave Simulations using Machine	
2	Learning: Kelvin-Helmholtz Instability and Mountain Wave Sources Driving	
3	Gravity Wave Breaking and Secondary Gravity Wave Generation	
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4		
5	Wenjun Dong <sup>1,2,3</sup> , David C. Fritts <sup>1,2</sup> , Thomas S. Lund <sup>2</sup> , Alan Z. Liu <sup>1</sup> , Han-Li Liu <sup>3</sup> ,	
6	and Jonathan Snively <sup>1</sup>	
7		
8	<sup>1</sup> Center for Space and Atmospheric Research (CSAR) and Department of Physical	
9	Sciences, Embry-Riddle Aeronautical University, Daytona Beach, FL, USA	
10	<sup>2</sup> Global Atmospheric Technologies and Sciences (GATS), Boulder, CO, USA	
11	<sup>3</sup> High Altitude Observatory, National Center for Atmospheric Research, Boulder, CO,	
12	USA	
13	Corresponding author: Alan Z. Liu	
14	Email: liuz2@erau.edu	
15	Key Words: Machine Learning, Kelvin-Helmholtz instability, Gravity Wave	
16	Key Points:	
17	1. A machine learning model (CAMNet) tailored for nonlinear gravity wave	
18	simulations is developed.	
19	2. CAMNet can achieve a several order-of-magnitude acceleration relative to	
20	physics-based model without sacrificing accuracy.	
21	3. CAMNet opens a new window to improve the parameterization of primary and	
22	secondary GWs in the global atmospheric models.	

#### 23 Abstract

Gravity waves (GWs) and their associated multi-scale dynamics are known to play 24 25 fundamental roles in energy and momentum transport and deposition processes throughout the atmosphere. We describe an initial, two-dimensional (2-D), machine 26 learning model - the Compressible Atmosphere Model Network (CAMNet) - intended 27 28 as a first step toward a more general, three-dimensional, highly-efficient, model for applications to nonlinear GW dynamics description. CAMNet employs a physics-29 informed neural operator to dramatically accelerate GW and secondary GW (SGW) 30 31 simulations applied to two GW sources to date. CAMNet is trained on high-resolution simulations by the state-of-the-art model Complex Geometry Compressible 32 33 Atmosphere Model (CGCAM). Two initial applications to a Kelvin-Helmholtz instability source and mountain wave generation, propagation, breaking, and SGW 34 35 generation in two wind environments are described here. Results show that CAMNet 36 can capture the key 2-D dynamics modeled by CGCAM with high precision. Spectral characteristics of primary and SGWs estimated by CAMNet agree well with those 37 from CGCAM. Our results show that CAMNet can achieve a several order-of-38 39 magnitude acceleration relative to CGCAM without sacrificing accuracy and suggests 40 a potential for machine learning to enable efficient and accurate descriptions of primary and secondary GWs in global atmospheric models. 41

#### 42 Plain Language Summary

Atmospheric gravity waves (GWs) are well described by the Navier-Stokes equations, 43 but solving these equations including small scale remains daunting, limited by the 44 very high computational cost of resolving the smallest spatial-temporal features in a 45 46 global context. To address this challenge, we developed a machine learning model called CAMNet. Our model demonstrates that neural networks can be trained on 47 high-resolution compressible atmospheric model data and then used to simulate 48 gravity wave evolution. Importantly, initial results show that using such trained 49 model can achieve computational savings of >1000 times compared to a physics-based 50

simulation while still achieve highly accurate results. These findings are exciting, as they suggest that CAMNet can overcome the limitations of current GW parameterizations and provide a promising avenue for studying the effects of subgrid-scale processes in atmospheric science and properly incorporating them in global models. The development of CAMNet opens up major new opportunities for improving effective model resolution, accuracy, and efficiency.

#### 57 1. Introduction

Gravity waves (GWs) play prominent roles throughout Earth's atmosphere. They are 58 generated at lower altitudes by various primary sources including airflow over 59 topography (i.e., mountain waves, MWs), convection, and jet streams. Additional 60 GWs are generated due to strong GW/mean-flow interactions described as "self-61 acceleration" (SA) dynamics and by resonant and off-resonant wave-wave 62 interactions that can take many forms. These diverse GWs play central roles in 63 64 Earth's atmospheric dynamics and climate by transporting energy, pseudomomentum, and constituents over depths extending into the thermosphere (Fritts & 65 Alexander, 2003). 66

Accounting for larger- and smaller-scale GW transports and influences remains a 67 challenging problem due to the complex physics involved and the need for high-68 69 resolution simulations to describe detailed responses where these are important. Traditional simulation methods, such as finite difference or finite volume methods, 70 can be computationally expensive, and quantifying sub-grid scale processes has been, 71 72 and remains, a major challenge. Multiple parameterization schemes spanning 40 73 years have aimed to account for GW pseudo-momentum deposition for various GW 74 sources, discretely or spectrally (i.e., linearly or nonlinearly), from the surface into the thermosphere (e.g., Lindzen, 1981; Holton, 1982; Palmer et al., 1986; Fritts and 75 Lu, 1993). More recent schemes have built on these earlier efforts and insights 76 (Warner and McIntyre, 1996; Hines, 1997; Alexander & Dunkerton, 1999; Yiğit et al., 77

2008; Eckermann et al., 2015; Amemiya & Sato, 2016; Gettleman et al., 2019; Miyoshi
& Yiğit., 2019; Ribstein et al., 2022).

Importantly, all these various schemes are based on simplified, often linear or weakly nonlinear, mathematical models and/or empirical relations that are significant approximations having limited quantitative predictive abilities. As such, they introduce significant model uncertainties and biases in predictions of middle and upper atmosphere responses (Pedatella et al., 2014). Additionally, parameter settings in these schemes may require adjustments for different models, model configurations, and/or model resolutions.

Parameterizations addressing GWs that are partially resolved (the "gray zone") that 87 88 maintain physical consistency between the resolved and parameterized dynamics are 89 promising (Vosper, 2015; Vosper et al., 2016), but can also be challenging (Liu, 2019). Such efforts attempt to represent the complex and highly nonlinear physics of GWs. 90 91 However, there remain many aspects of GW dynamics, e.g., SA dynamics, local 92 instabilities and breaking, multi-scale interactions, and secondary GW (SGW) 93 generation that become increasingly important at increasing altitudes, but that 94 cannot be addressed by linear theory or existing GW parameterization schemes.

95 The recent boom in hardware and software developments relevant to machine
96 learning (ML) has motivated some efforts to examine the possible benefits that ML
97 can bring to GW parameterization (e.g., Chantry et al., 2021; Espinosa et al., 2022).
98 ML methods offer several potential advantages over traditional parameterization
99 schemes. These include the following:

100 1) ML applications can learn complex, nonlinear relationships directly from 101 data, without the need for pre-determined equations or assumptions. This 102 makes ML methods well-suited for problems exhibiting highly nonlinear 103 dynamics;

104 2) ML can identify complex patterns and relationships in the data that may be
105 difficult to discern using traditional methods. This can lead to improved
106 accuracy and transferability of the resulting flow descriptions; and

3) ML has the potential to dramatically reduce the computational cost
associated with descriptions of nonlinear GW dynamics by replacing
traditional parameterizations with highly-efficient, data-driven models.

Such methods have been shown to yield significant benefits in several applications to 110 111 date. Chantry et al. (2021) trained a neural network on an upgraded version of an 112 existing parameterization scheme that yielded improved results describing GW drag in a numerical weather prediction (NWP) system. In another study, Espinosa et al. 113 114 (2022) developed an artificial neural network to emulate the pseudo-momentum forcing described in a traditional GW parameterization in an idealized climate model. 115 116 By coupling the climate model with their ML-based GW parameterization, they were 117 able to accurately reproduce the quasi-biennial oscillation, a well-known atmospheric phenomenon. However, these ML-based GW parameterizations rely on traditional 118 GW parameterizations, hence inherit their assumptions and simplifications. They 119 also remain unable to represent the true, highly nonlinear GW dynamics. 120

121 As mentioned earlier, the dynamics of GWs are governed by the Navier-Stokes equations. In recent years, several ML-based solvers for partial differential equations 122 123 (PDEs) have been proposed to approximate or improve various numerical methods. 124 The most explored of these can be divided into two categories: physics-informed 125 neural networks (PINN, e.g., Maziar et al., 2019; Wandel et al., 2022) and neural 126 operators (NOs, e.g., Lu et al., 2019; Li et al., 2020; Xiong et al., 2023). PINN uses a 127 neural network as the solution function and optimizes a loss function to minimize violation of the given equation. However, it experiences difficulties in propagating 128 information from initial or boundary conditions to unseen parts of the interior and to 129 future times. NOs are better suited for solving PDEs and have been successfully used 130 in flow prediction (e.g., Lu et al., 2019; Li et al., 2020; Xiong et al., 2023). However, 131

they require large volumes of simulation data. Recently, physics-informed neural operators, e.g., both the physics-informed Deep Operator Network proposed by Goswami et al(2022) and physics-informed Fourier neural operator proposed by Li et al (2022) employ both data and physics losses on operator learning to overcome the shortcomings of purely PINN or data-driven learning. Our ML approach is also based on this physics-informed neural operator.

138 While MLs offer potential advantages for GW simulation, there are also challenges 139 that must be addressed. The primary challenge is the limited availability of highquality data for training and validation. To address this need, we utilize the Complex 140 Geometry Compressible Atmospheric Model (CGCAM), a finite volume model that 141 142 has been used extensively to study GW dynamics and their instabilities in the Mesosphere and Lower Thermosphere (MLT) at very high resolution, see, e.g., Dong 143 144 et al., 2020, 2021,2022,2023; Fritts et al., 2020, 2021,2022a, 2022b; Lund et al., 2021). CGCAM is capable of capturing highly nonlinear GW and GW-related dynamics 145 yielding high-fidelity GW training datasets for ML-based approaches. 146

147 Inspired by previous research in this field, this study investigates the application of 148 ML algorithms for simulating GW dynamics. Our focus is on the potential of ML to 149 improve the efficiency of GW simulations while capturing their highly nonlinear dynamics with high fidelity, specifically including instabilities, breaking, and SGW 150 generation. Our approach employs the Compressible Atmosphere Model Network 151 152 (CAMNet) model based on CGCAM that solves the compressible Navier-Stokes equations in the Complex Geometry Compressible Atmosphere Model (CGCAM). To 153 explore the performance of CAMNet, we train the model using single-channel inputs 154 instead of multi-channel inputs as explored by Pathak et al (2022), with a focus on 155 156 improving training efficiency. We expect this study to provide valuable insights into the potential benefits of ML in modeling GW and instability dynamics and to inform 157 future research in this area. When optimized, this method will be extended to three-158 dimensions in order to achieve much more efficient descriptions of GW effects in 159 160 global models while achieving sufficient accuracy.

## 161 2. Method

#### 162 2.1 Complex Geometry Compressible Atmospheric Model (CGCAM)

163 The Complex Geometry Compressible Atmosphere Model (CGCAM) solves the
164 two/three-dimensional (2/3-D) nonlinear and compressible Navier-Stokes equations
165 written in strong conservation law (flux) form as follows:

166 
$$\frac{\partial \rho}{\partial t} + \frac{\partial (\rho u_j)}{\partial x_j} = 0$$
(1)

167 
$$\frac{\partial(\rho u_i)}{\partial t} + \frac{\partial(\rho u_i u_j)}{\partial x_j} = -\frac{\partial p}{\partial x_i} - \rho g \delta_{i3} + \frac{\partial \sigma_{ij}}{\partial x_j}$$
(2)

168 
$$\frac{\partial(\rho E)}{\partial t} + \frac{\partial\left[(\rho E + p)u_j\right]}{\partial x_j} = -\rho g u_3 + \frac{\partial\left(u_i \sigma_{ij}\right)}{\partial x_j} - \frac{\partial q_j}{\partial x_j}$$
(3)

169 where  $\sigma_{ij}$  and  $q_j$  are the viscous stress and thermal conduction, respectively, defined 170 as

171 
$$\sigma_{ij} = \mu \left[ \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \frac{2}{3} \left( \frac{\partial u_k}{\partial x_k} \right) \delta_{ij} \right]$$
(4)

$$q_j = -\kappa \frac{\partial T}{\partial x_j} \tag{5}$$

and where  $\mu$  is the dynamic viscosity,  $\kappa$  is the thermal conductivity,  $\delta_{ij}$  is the Kronecker delta,  $\rho$  is density, and g is the gravitational acceleration.  $\mu$  and  $\kappa$  depend on the temperature through Sutherland's Law (White, 1974).

176 The solution variables are  $\rho$ , the momentum per unit volume,  $\rho u_i \text{ or } (\rho u, \rho v, \rho w)$ , and 177 the total energy  $E = e + u_k u_k/2 = c_v T + u_k u_k/2$ , with velocity components 178  $(u_i, u_j, u_k)$  along (x, y, z). Also  $c_v = \frac{R}{\gamma - 1}$  is the specific heat at constant volume and *T* is 179 the temperature. The compressible equation set is discretized using a second-order finite-volume scheme identical to the method discussed by Felten & Lund (2006).
Time advancement is achieved via a third-order accurate Runge-Kutta scheme.
Additional details for CGCAM are provided by Dong et al. (2020) and Lund et al.
(2020). CGCAM has been successfully used in various studies on GW generation,
breaking, and SGW generation (e.g., Dong et al., 2020,2021,2022; Fritts et al.,
2020;2021; 2022a; 2022b; Lund et al., 2021).

#### 186 2.2 Compressible Atmospheric Model Network (CAMNet)

187 CAMNet is a hybrid machine learning model that combines data-driven and physicsinformed approaches. It is based on the Adaptive Fourier Neural Operator (AFNO) 188 proposed by Guibas et al., 2021. AFNO is a Fourier transform-based token-mixing 189 scheme with a vision transformer backbone (Dosoviskiv et al., 2020). AFNO is based 190 191 on the Fourier neural operator (FNO) that learns in a resolution-invariant manner 192 and has shown success in modeling challenging partial differential equations such as 193 fluid dynamics (Li et al., 2020). The Fourier architecture of AFNO applies a fast 194 Fourier Transform (FFT) to the data and applies its fully connected layers in Fourier 195 space before performing an inverse FFT back to real space. Moreover, the Fourier 196 architecture has been demonstrated the ability to perform zero-shot super-resolution, 197 predicting on higher-resolution data having only seen low resolution data. The introduction of vision transformer enables it to model long-range dependencies well 198 and yields a state-of-the-art high-resolution model that resolves fine-grained features 199 200 and scales well with resolution and size of dataset. AFNO enables training high-201 fidelity data-driven models as truly unprecedented resolution (Pathak et al., 2022).

The power of AFNO stems from its ability to combine linear integral operators, implemented through the Fourier transform, with non-linear activation functions, enabling it to learn highly non-linear operators. This is similar to standard Multi-Layer Perceptron (MLP) and Convolutional Neural Network (CNN), where linear multiplications are combined with non-linear activations to learn highly non-linear functions. Although AFNO truncates higher frequency modes in the Fourier layer, Li 208 et al. (2020) argue that the entire operator can still approximate functions with the 209 full frequency range, due to the function being represented in a high-dimensional channel space. The non-linear decoder network then recovers the higher frequency 210 modes when projecting back to the desired dimension. In our case of predicting multi-211 scale GW dynamics, the Fourier layer truncation of high-frequency information 212 213 resulted in poor small-scale structure prediction. To address this, we added a convolutional layer, which is able to amplify high-frequency components and 214 215 complement the information truncated by the Fourier layer.

Additionally, CAMNet further extends the AFNO architecture by incorporating 216 physics information from the Navier-Stokes equations, which govern the GW 217 218 dynamics. These equations are used to create a loss function that captures the violation of these laws, and Fourier derivatives (Li et al., 2021) are used to compute 219 the derivatives for the physics constraints, as the automatic differentiation in 220 PyTorch is very memory intensive for this type of architecture. The physical 221 constraints reduce the demand for training datasets and improve the generalization 222 223 and physical validity of CAMNet learning compared to purely data-driven methods.



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225 Figure 1. CAMNet architecture that utilizes the modified Adaptive Fourier Neural Operator (AFNO) and follows a patch-based approach. The input frame is divided 226 into a  $h \times w$  grid of patches, each of size  $p \times p$ , and encoded in a higher dimensional 227 228 space with position embedding is added to form a sequence of tokens. These tokens 229 are then mixed spatially using AFNO, which is repeated for L layers, and then a 230 decoder reconstructs the patches for the next frame. A single AFNO layer is composed of multiple heads for parallel processing. Parts 2-4 of the layer applies the 231 Fourier transform  $\mathcal{F}$  to the input, followed by a linear transform  $\mathcal{R}$  that acts on the 232 233 lower-frequency Fourier modes and filters out the higher-frequency modes, this is then followed by an inverse Fourier transform  $\mathcal{F}^{-1}$ . Part 5 stands for a convolutional 234 235 layer that is used to capture the higher-frequency modes that are missed by Part 2-4. Both AFNO outputs  $v_L$  and reconstructed U are functions, and their derivatives 236  $Dv_L$  and DU can be computed at any query points x and z. The solutions are 237 constrained by the Navier-Stokes equations that govern GW dynamics. For more 238 239 information, refer to the corresponding texts.

The CAMNet architecture is shown in **Figure 1**. The CAMNet is currently trained on horizontal wind U, so CAMNet's input is U(x, z, t). The CAMNet architecture consists of 7 parts. Below, we present a detailed computational implementation of each part. Similar to the iterative update strategy of each Fourier layer in FNO (Li et al., 2020), the improved iterative update strategy of each Fourier layer in CAMNet can be expressed as follows:

246 
$$v_1 = \stackrel{Part \ 1}{\widetilde{\mathcal{P}}} U(x, z, t) \tag{6}$$

247 
$$\boldsymbol{v}_{l+1} = \sigma \begin{pmatrix} part \ 5 & part \ 2 & part \ 3 & part \ 4 \\ \widehat{\mathcal{C}} & \boldsymbol{v}_l + \widetilde{\mathcal{F}^{-1}} & \widehat{\mathcal{R}} & \widehat{\mathcal{F}} & \boldsymbol{v}_l \end{pmatrix}, for \ l = 1, 2, \cdots, L$$
(7)

248 
$$U(x, z, t+1) = \overset{part 6}{\tilde{Q}} v_L$$
(8)

where  $\mathcal{P}$  and  $\mathcal{Q}$  are encoder and decoder that are realized by two neural networks that projects U(x, z, t) to hidden representation  $v_l$  and projects the representation back to the solution U(x, z, t + 1).  $\sigma$  is a nonlinear activation function. The additional term  $\mathcal{C}$ is a convolutional layer that acts on  $v_l$ .  $v_l$  denotes the output of the *l*-th Fourier layer of AFNO.  $\mathcal{F}$  denotes the Fourier transform that acts on  $v_l$ ,  $\mathcal{R}$  is a linear transform layer that acts on  $\mathcal{F}(v_l)$  to handle its low-frequency modes.  $\mathcal{F}^{-1}$  is an inverse Fourier transform that acts on  $\mathcal{R}(\mathcal{F}(v_l))$ . The details of each part are provided below:

**Part 1: Encoder**  $\mathcal{P}$  The encoder is implemented using the token embedding layer in the Vision Transformer architecture proposed by Dosovitskiy et al. (2020). This layer applies a linear projection to each patch to obtain a fixed-sized vector, which is then concatenated with positional embeddings representing the spatial location of the patch. The resulting sequence of vectors serves as the input to the sequent neural network layers.

## 262 Part 2-4: Fourier Transform $\mathcal{F}$ , Linear Transform $\mathcal{R}$ , and Inverse Fourier Transform 263 $\mathcal{F}^{-1}$

264 As stated in Li et al., 2020, since the inputs and outputs of partial differential 265 equations (PDEs, such as Navier-Stokes equations) are continuous functions, it is more efficient to represent them in Fourier space and perform global convolution. 266 267 This is due to the quasi-linear computational complexity and global properties of 268 Fourier transform, making it a more efficient approach. The convolution in the spatial 269 domain is equivalent to the pointwise multiplication in the Fourier domain. To 270 capture global features in input data, a Fourier transform is first applied to the inputs, followed by a Linear Transform  $\mathcal{R}$  that acts on the lower-frequency Fourier modes by 271 assigning weights to them. These weights will be updated during the training. Finally, 272 273 an inverse Fourier transform is performed to obtain the output.

274 Part 5: High-frequency Information Compensation C In each Fourier layer, we utilize
275 a convolutional layer to extract high-frequency information because it can amplify

high-frequency components. Therefore, we train a convolutional layer c on the outputs of Part 1 to extract their high-frequency information. As a complement to Parts 2-4, the convolutional layer enables the forward prediction of high-frequency information.

Part 6: Decoder *Q* Given two future states independently predicted by Parts 2-4 and
Part 5, we combine them and train a non-linear decoder using a multi-Layer
perceptron layer with a tanh activation function to transform the AFNO outputs back
into *U*.

#### 284 Part 7: Physics Informed Loss $\mathcal{L}_p$

As CAMNet is currently trained on horizontal wind U, we utilize only the momentum flux equation (1) for the physics-informed part, assuming a constant density  $\rho$  over time. Thus, equation (1) can be simplified to

288 
$$\frac{\partial(\rho U)}{\partial x} + \frac{\partial(\rho W)}{\partial z} = 0$$
(9)

In equation (9), the vertical wind W is obtained from CGCAM simulations at each prediction time step, and  $\rho$  is set to its initial values. The only physics-informed variable to be calculated in CAMNet is the derivative of U with respect to x. To enable multi-variable predictions in the future, the physics-informed part will need to involve the complete Navier-Stokes equations.

294 The loss function of CAMNet for optimizing equations (6), (7), and (8) is defined as

295 
$$\mathcal{L} = \alpha \mathcal{L}_{data} + \beta \mathcal{L}_{p} = \alpha \left| U - \widehat{U} \right| + \beta \left| \frac{\partial(\rho U)}{\partial x} + \frac{\partial(\rho W)}{\partial z} \right|$$
(10)

296 where  $\alpha$  and  $\beta$  control the weights of data-driven and physics-informed part in loss 297 functions, respectively.

**Table 1**: Key model parameters used in CAMNet model and training.

Hyperparameter	Value
Batch Size	32
Learning rate	$5 \times 10^{-4}$
Learning rate schedule	Cosine
Patch size (Case 1/2)	$10 \times 10/8 \times 8$
Number of AFNO layers	8
Heads number	8
Heads depth	6
AFNO embedding dimension	768
Activation function	GeLU
Dropout	0

299 CAMNet is highly optimized so that it can be trained efficiently on massively parallel 300 GPU resources. The initial application of CAMNet has shown that CAMNet can 301 achieve order-of-magnitude speedup over numerical model CGCAM. We performed 302 two cases to evaluate the performance of CAMNet (see "Result" section), and the 303 relevant model parameters are presented in **Table 1**.

## 304 3. Results

This section describes our initial efforts using CAMNet to accelerate simulations of 305 306 GWs arising from two very different sources constrained to a two-dimensional (2-D) 307 domain. Case 1 describes the generation of initial GWs by large-scale, shear-induced Kelvin-Helmholtz Instability (KHI) and the successive generation of SGWs at much 308 larger scales that readily propagate to much higher altitudes. Case 2 describes 309 mountain waves arising from flow over idealized terrain, their attainment of large 310 amplitudes, breaking, and generation of SGWs that likewise attain very high 311 312 altitudes. CAMNet wind fields and spectra are compared with high-resolution 2-D 313 CGCAM simulations in both cases.

#### 314 3.1 Case 1: Gravity Waves emitted from Kelvin–Helmholtz Instability

We explore here the potential of CAMNet for modeling GWs emitted from Kelvin-Helmholtz Instability (KHI) described by Dong et al.(2023). We use CGCAM to generate training and testing data for CAMNet. The initial background winds (see Figure 2) are specified as

319 
$$U(z) = U_0 cos\left[\frac{\pi(z-z_0)}{15 \ km}\right] tanh\left(\frac{z-z_0}{h}\right)$$
(11)

320 CGCAM simulations are performed for a computational domain having dimensions 180 km  $\times$  180 km (x, z) with resolutions of 50 m at the shear center, with exponential 321 mesh stretching approaching the upper and lower boundaries to reduce 322 computational demands. Periodic boundary conditions are used at the lateral 323 324 boundaries. An isothermal no-stress wall condition is used at the lower boundary, 325 and a characteristic radiation condition is used at the upper boundary. The vertical boundary conditions are supplemented with sponge layers having 20-km depths to 326 327 further ensure no reflected GWs. After excluding irrelevant data in the sponge layers, the variable U are stored on a grid of dimension of  $2000 \times 1000$ . 328



Figure 2: Initial conditions for generating training and testing data (left) and an
example of initial U in the simulation domain (right) for the KHI cases.

Given the initial conditions, CAMNet is required to simulate the future states of 332 333 variable *U* at  $t \in \{1, 2, 3, \dots, 50\}$  min for a suite of initial conditions. A total of 200 cases are generated by varying  $U_0$  and *h* in Equation 11, and the corresponding outputs of 334 CGCAM serves as the true reference solutions for each case. The CGCAM simulations 335 for each case were run for 50 minutes at an interval of 1 minute. The 200 CGCAM 336 337 cases were then split into a training set of 180 cases and a testing set of 20 cases. The training set is used to train CAMNet, and the testing set is used to evaluate the 338 model's performance. Keeping the testing set separate from the training set is crucial 339 340 to obtain an unbiased estimate of the model performance. All samples have a grid of 341  $2000 \times 1000$ . The CAMNet training is implemented in a multi-GPU environment 342 with 4 V100 GPUs. Convergence is observed after approximately 250 epochs during 343 the training process.



344

345

Figure 3. KHI and GWs predicted by CGCAM and CAMNet.

Figure 3 displays the variable *U* employed for both the CAMNet and CGCAM simulations during the model test, and we observe a high level of consistency between two models. The ML model CAMNet can capture small- and large-scale structures qualitatively, with clear evidence of KHI and KHI-radiated GWs seen in both CGCAM and CAMNet. The initial strong shear produces deep and broad KH billows

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that break down after ~30 min, leading to the emergence of small-amplitude GWs above the KHI altitude. By 50 min, increasing GWs and apparent SGWs are seen to propagate with group velocities along the iso-phase lines that extend away from the KHI shear layer, causing GWs to propagate to higher altitudes and achieve larger amplitudes. Both models suggest that GWs are continuously emitted from the KHI and turbulence dynamics, with well-defined spatial structures having orientations and spatial scales that agree closely between the CGCAM and CAMNet fields.



358

Figure 4. Case 1 spectral characteristics predicted by CGCAM (top) and CAMNet
(bottom) at 70-80 km (left) and 90-120 km (right) altitude ranges at 30, 40, and 50
min model simulation times (see time labels at lower right).

The spectral properties of GWs can reveal important details about their sources, such as the altitude and vertical extent of the source region, and the dominant wavelengths and frequencies of the GWs generated. This information can be used to improve parameterizations of GW sources in global climate models, which are critical 366 for predicting the response of the atmosphere to changing climate conditions. The spectral structures of horizontal wind disturbance u' are calculated and found to 367 exhibit high consistency between CGCAM and CAMNet. The spectra of  $u^{\prime 2}$  computed 368 at the KHI region and higher altitudes at t = 30, 40, and 50 min are displayed in 369 370 Figure 4. The spectra reflect the characteristics of KHI dynamics. In the KHI region, 371 the onset of strong KHI and 2-D turbulence yields spectral slopes approaching -5/3 corresponding to wavenumbers of ~1-6 rad/km. Spectral amplitudes fall sharply 372 beyond wavenumbers of ~20 rad/km and exhibit steeper slopes approaching -7 in the 373 374 viscous range. At higher altitudes, the -5/3 slope corresponds to wavenumbers of ~0.8-375 3 rad/km and -7 slopes correspond to wavenumbers of 3-10 rad/km. These spectra 376 suggest that small-scale structures discussed above are well resolved at these times.

These results suggest that CAMNet has promising potential as an alternative to CGCAM for simulating KHI and KHI-radiated GWs. Notably, the time cost of a single KHI case simulation using CAMNet was approximately 0.8 seconds on a single A100 (80GB). This represents a significant acceleration (by > 2000) compared with the 30 minutes needed by CGCAM when using 36 CPU cores.





Figure 5. Initial mean winds U(left) and temperature T(right) as the inputs of
 CGCAM for generating training and testing data for MW cases.

386 Our intent in Case 2 was to explore CAMNet capabilities for modeling MW generation, propagation, breaking, and their radiation of SGWs. MW breaking is one of the 387 strongest sources of SGWs (Lund et al., 2020). As in Case 1, CAMNet training 388 employs CGCAM simulation data. The CGCAM simulations cover a computational 389 domain extending  $700 \times 220$  km (x and z) at a resolution of 1 and 0.5 km in x and z. 390 391 As in Case 1, the lateral boundary is periodic. At the lower boundary, a Gaussian 392 terrain of peak height 4 km and half-width of 30 km is used, and a characteristic radiation condition is used at the upper boundary. Sponge layers of 20 km and 50 km 393 394 are added to the vertical and lateral boundaries, respectively, to ensure absorption of 395 outgoing GWs. The variable U are stored on a grid having dimensions of  $600 \times 400$ 396 after irrelevant values in the sponge layers are excluded.



- **Figure 6**. MW evolutions predicted by CGCAM (left) and CAMNet (right). The
- initial condition is a horizontal wind of 25 m/s at the surface flowing over a
- 400 Gaussian Mountain with a height of 4 km and a half width of 30 km (see Figure 4).

Given the samples of initial conditions, CAMNet is trained to reproduce the future 401 states of variable U at  $t \in \{140, 145, 150, \dots, 220\}$  min. Note that we start from t = 140402 mins to avoid CAMNet being trained with non-physical data produced by CGCAM at 403 404 early simulation times. A total of 200 cases are generated by varying the initial wind field, with a random wind field randomly extracted from the HWM14 at  $30^{\circ}S$ ,  $70^{\circ}W$ 405 (Andes Lidar Observatory) at 00:00 on 200 days among 365 days. The initial 406 temperature field is simplified as used in Dong et al. (2020). Winds at lower altitudes 407 408 from HWM14 are consistently lower than actual observations, and thus a correction is needed to enable simulation of MW generation. To account for this discrepancy, we 409 randomly assign wind values ranging from 0.30 m/s at these lower altitudes to 410 411 facilitate the occurrence of MWs. The initial fields were assumed to be uniform over 412 the domain. The initial wind and temperature fields are shown in Figure 5. The 413 corresponding output of CGCAM serves as the true reference solution for each case. 414 CGCAM simulations for each case were run for 220 minutes at an interval of 5 415 minutes. The CGCAM simulations were then split into a training set of 180 cases and 416 a testing set of 20 cases. All samples have a grid of 600 × 400. The training of CAMNet 417 is implemented in a multi-GPU environment with 4 V100. Convergence is observed after approximately 320 epochs during training process. 418

The CAMNet model achieves excellent skill in modeling MW generation, propagation, 419 breaking, and SGW generation. As an illustrative example, we choose a case with a 420 uniform initial background wind of 25 m/s. We begin with an overview of the major 421 features of the MW and SGW evolution from t = 150-220 min. Figure 6 shows U at 422 423 t = 150, 180, 200, 220 min generated by CAMNet and the corresponding CGCAM results at these times. Considering CAMNet results first, the earliest responses at 424 t = 150 min reveal MW generation at lower altitudes and their extension into the 425 426 MLT. At t = 180 min, initial SA dynamics and instabilities are seen at lower altitudes.

427 At t = 200 min, there is evidence for strong SGW excitation in the MW breaking 428 regions. The MW field and its associated instabilities and SGWs continue to intensify 429 to t = 200 min. The CAMNet results approximate the CGCAM ground truth 430 remarkably well over 220 min. Additionally, high consistency is found between 431 CGCAM and CAMNet in the u' spectra, which are shown in the first and second rows 432 of Figure 8, respectively.





Figure 7. Same as Figure 6, but for a tidal wind background.



435

Figure 8. Case 2 spectral characteristics predicted by CGCAM (first and third rows)
and CAMNet (second and fourth rows) at 70-80 km (left), 100-120 km (middle) and
90-120 km (right) at 150,180,200, and 220 min (see time labels at lower right).

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As a further test of the generalization ability of the trained CAMNet model, we chose
an initial wind field that includes a tidal wind field (represented by grey lines in
Figure 7) that was not part of the 200 cases used for training and testing. CAMNet
results were compared to those of CGCAM, and also showed high consistency in
capturing the MW dynamics (see Figures 7 and 8), including the following:

- 444 1. Major MW responses, including MW generation, propagation, breaking, local
  445 instabilities and dissipation, and SGW generation;
- 446 2. strong MW breaking, instabilities, and 2-D turbulence dynamics;
- 3. SGWs modulated by tidal winds and having large scales and large influences
  extending into the thermosphere. Responses include refraction by tidal winds,
  and reduced SGW responses at higher altitudes relative to the case of no tidal
  wind (as shown in Figures 5 and 6); and
- 4. CAMNet exhibits highly consistent MW spectral characteristics with CGCAM452 (see the third and fourth rows in Figure 7).

453 This case demonstrates that CAMNet exhibits good generalization abilities.

454 CAMNet demonstrates promising potential as a competitive alternative to CGCAM
455 for simulating MW generation, propagation, and breaking. A single MW case
456 simulation using CAMNet takes approximately 0.5 seconds on a single A100 (80GB),
457 in this case ~4000 times faster than the corresponding CGCAM simulation using 36
458 CPU cores, which took around 40 minutes.

#### 459 4. Discussion

GWs play significant roles in the transport of energy and pseudo-momentum through Earth's atmosphere. However, current numerical weather prediction and climate models lack the resolution needed to describe the smaller-scale GWs and SGWs accounting for the large majority of GW energy and pseudo-momentum fluxes to higher altitudes. Thus, parameterizations are employed to represent unresolved GW influences in most global atmospheric models. These typically have three primary 466 components: (1) specification of GWs at source levels, (2) GWs propagation with467 altitude, and (3) GW dissipation and further parameterized body forcing and mixing.

468 The primary function of GW parameterizations as currently applied in global models is to compute the wave-driven force on the mean flow, where the mean flow refers to 469 the grid-box mean, and the waves are intended to represent sub-grid unresolved GW 470 anomalies. These parameterizations treat GWs as linear, hydrostatic, vertically 471 472 propagating waves in a steady ambient environment with Boussinesg governing 473 equations. However, these assumptions and simplifications limit the representation of many observed GW characteristics, which are physically well understood (e.g., 474 Erving et al., 2006, Hertzog et al., 2012; Eckermann et al., 2015; Stephan et al., 2016). 475 476 Some ML-based GW parameterization schemes have been proposed to enhance the 477 accuracy of GW parameterization (e.g., Chantry et al., 2021; Espinosa et al., 2022). These ML-based GW parameterization schemes rely on traditional GW 478 479 parameterizations and hence are limited by the assumptions and simplifications 480 inherent in them. CAMNet offers several advantages over traditional and previous 481 ML-based GW parameterization schemes. The training process of CAMNet does not 482 rely on any existing GW parameterization schemes, thus it is not limited by their 483 assumptions and simplifications. CAMNet is trained with high-resolution simulation 484 data from the state-of-the-art atmospheric model CGCAM, which are accurate 485 numerical solutions of the Navier-Stokes equations. Well-trained CAMNet is capable of resolving multi-scale and highly nonlinear GW dynamics, such as instability and 486 turbulence dynamics, self-acceleration, GW breaking, and SGW generations, at much 487 488 faster speed. To the best of our knowledge, CAMNet is the first ML-based approach 489 that can directly simulate highly nonlinear GW dynamics.

490 To further enhance the application of CAMNet's exceptional GW simulation 491 capabilities, CAMNet will be optimized and extended 1) from single-variable 492 simulation to multi-variable simulation to better capture the complex nonlinear 493 interactions among various GW variables; and 2) from 2D to 3D to more accurately

494 model the horizontal evolution characteristics of GWs, which are only fully displayed495 in 3D cases.

#### 496 5. Summary

497 In this paper, we developed a machine learning model solving the compressible 498 Navier-Stokes equations in our Complex Geometry Compressible Atmosphere Model (CGCAM) named CAMNet. CAMNet is a hybrid machine learning model that 499 500 combines data-driven and physics-informed approaches. It is based on the Adaptive 501 Fourier Neural Operator (AFNO) proposed by Guibas et al., 2021, with modifications tailored to our simulations. The main improvements include: 1) the addition of 502 503 convolutional layer branch to compensate for high-frequency components truncated 504 by the Fourier layers, making the model more robust in resolving multi-scale 505 dynamics, and 2) the incorporation of physical information from the Navier-Stokes equation in CGCAM. The CAMNet feedback neural network utilizes a loss function 506 507 that combines both the physical information in Navier-Stokes equations and the data loss from CGCAM simulations. This approach reduces the need for extensive training 508 data and improves the model generalization ability. 509

510 We evaluated the performance of CAMNet with two test cases: the first one described 511 the KHI and the associated GW radiation as explored in Dong et al., 2023, and the 512 second one addressed the generation, propagation, and breaking of MWs, which are 513 one of the most important GW components. Both cases involve small-scale instability 514 and turbulence dynamics, as well as larger-scale GWs. Our main findings from the 515 simulations are as follows:

- CAMNet shows excellent skill on simulating the formation and intensification
   of KHI, and KHI-radiated GWs. Those results qualitatively match the ground
   truth remarkably well over a period of 50 mins.
- 519 2) CAMNet has excellent skill on the simulations of MW generation, propagation,
  520 and breaking. CAMNet captures the major MW responses, including MW

generation, propagation into MLT, strong MW breaking via vortex ring
formation, strong instability and turbulence dynamics, intense SGW
generation, and strong modulation by tidal wind.

524 3) CAMNet can be trained on high-resolution CGCAM simulations under various
525 conditions and has the potential to significantly accelerate GW simulations
526 while maintaining high accuracy.

527 Accounting for the GW forcing in global atmospheric models is challenging due to their limited model resolutions. A well-trained CAMNet can produce the simulations 528 orders of magnitude faster than CGCAM without any noticeable accuracy loss. This 529 has two important implications. Firstly, high-resolution GW simulations can be 530 531 generated within seconds, thus enabling estimation of well-calibrated and constrained uncertainties regarding unresolved GW scales with higher confidence 532 compared to current global models that have severely simplified GW 533 534 parameterization schemes due to computational cost. Secondly, CAMNet is suitable for rapidly testing hypotheses about mechanisms of GW forcing and their 535 536 predictability. Moreover, there is potential to develop a software library of welltrained CAMNet models to be applied to a broad range of conditions. The well-trained 537 models have the potential to become a viable alternative to current GW 538 parameterizations in global models. 539

#### 540 Data availability

541 Outputs from model simulations used in this study are archived on the ERAU High542 Performance Computing System Vega and can be made available upon publication.

#### 543 Code availability

544 The CGCAM code utilized in this study is proprietary and can only be accessed 545 through collaboration. The original AFNO code is available at 546 https://github.com/lonestar686/AdaptiveFourierNeuralOperator.

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